

Bimodules and non-commutativity in fundamental physics

by John Barrett

C vs NC

space $\longleftrightarrow$ commutative $\ast$ -algebra
vector bundle $\longleftrightarrow$ module over algebra

"NC space" $\longleftrightarrow$ NC $\ast$ -algebra $\mathcal{A}$

"NC vector bundle" $\longleftrightarrow$ bimodule
over algebra $\mathcal{H}$

Real structure

$$\mathcal{J}: \mathcal{H} \rightarrow \mathcal{H} \quad \text{antilinear}$$

$$\mathcal{J}^2 = \pm 1$$

$$\mathcal{J}(a\psi b) = b^*(\mathcal{J}\psi)a^*$$

Examples

- Standard model internal space
- Fuzzy coadjoint orbits
- Fuzzy torus : spin structure, π , etc.

SM algebra

$$A \ni a = \left(\begin{array}{c|c|c} \mathcal{L} & \cdot & \cdot \\ \hline \cdot & \begin{array}{c} \lambda \cdot \\ \cdot \bar{\lambda} \end{array} & \cdot \\ \hline \cdot & \cdot & \begin{array}{c|c} \lambda & \cdot \\ \hline \cdot & m \end{array} \end{array} \right)$$

$$\mathcal{L} = \begin{pmatrix} \alpha & \beta \\ -\bar{\beta} & \bar{\alpha} \end{pmatrix} \in \mathbb{H}$$

$$\lambda \in \mathbb{C}$$

$$m \in M_3(\mathbb{C})$$

Gauge group: $G = \{a \mid aa^* = 1, \det a = 1\}$

$$\leadsto \mathcal{L} \in \mathrm{SU}(2), \lambda \in \mathrm{U}(1)$$

$$m = \mu s, s \in \mathrm{SU}(3), \lambda \mu^3 = 1$$

SM internal space

$$\mathcal{H} \ni \Psi = \begin{pmatrix} & & & & \begin{matrix} \nu_L & u_L & u_L & u_L \\ e_L & d_L & d_L & d_L \\ \nu_R & u_R & u_R & u_R \\ e_R & d_R & d_R & d_R \end{matrix} \\ & 0 & & & \\ \hline \begin{matrix} \bar{\nu}_L & \bar{e}_L & \bar{\nu}_R & \bar{e}_R \\ \bar{u}_L & \bar{d}_L & \bar{u}_R & \bar{d}_R \\ \bar{u}_L & \bar{d}_L & \bar{u}_R & \bar{d}_R \\ \bar{u}_L & \bar{d}_L & \bar{u}_R & \bar{d}_R \end{matrix} & & & 0 \end{pmatrix}$$

Bimodule :
 $a \psi a'$
 matrix x .

$\otimes \mathbb{C}^3$
 (generations)

Action of $a \in G$: $\Psi \mapsto a \Psi a^*$ "adjoint"
 $\leadsto$ SM charges

Coadjoint orbits

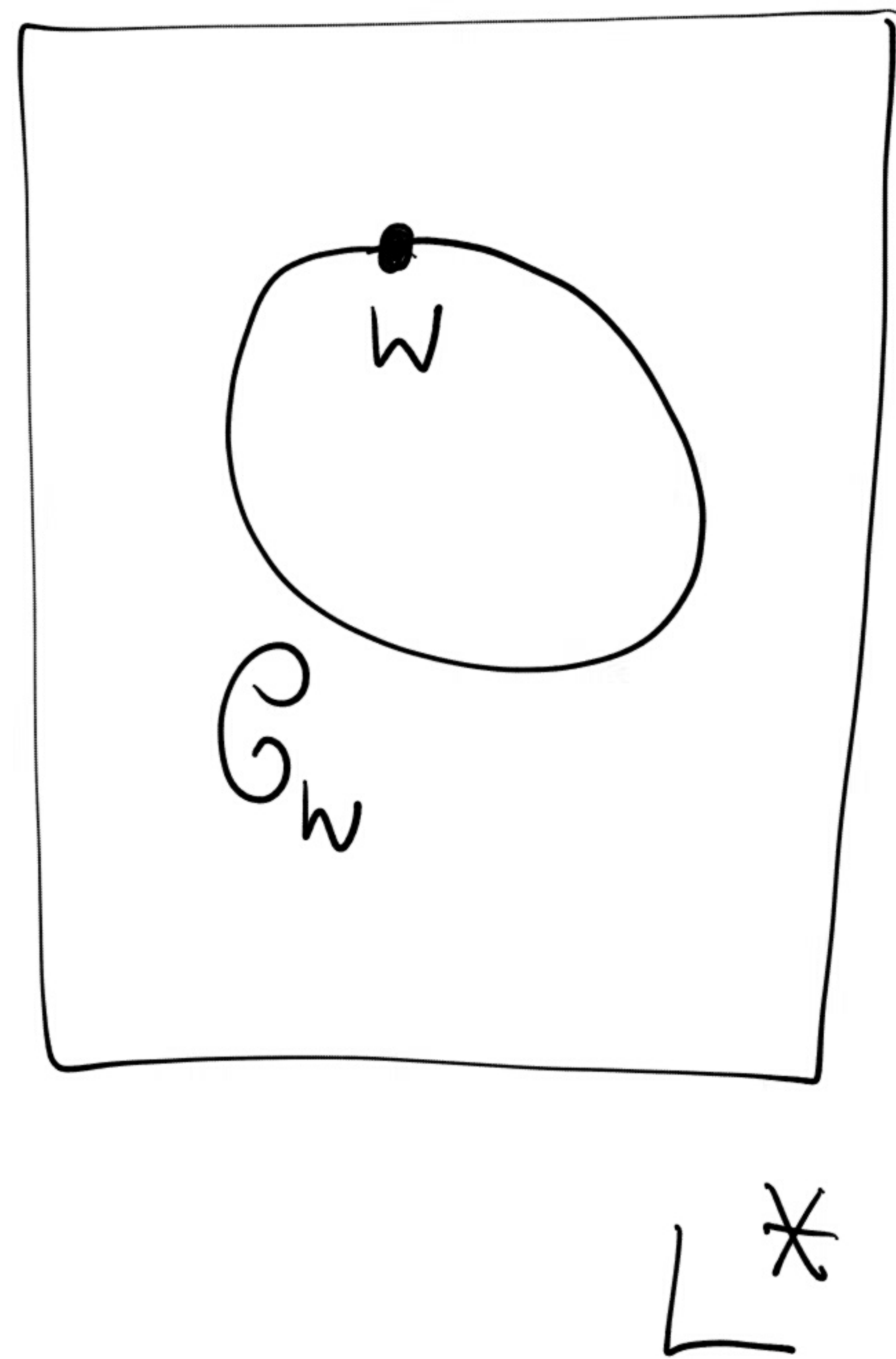
G : Lie group

L : its Lie algebra

$w \in L^*$

$$\mathcal{O}_w = G \{w\} \cong G/H$$

G semisimple $\leadsto \mathcal{O}_w$ Kähler.



Functions on coadjoint orbit

$L^2(\mathcal{O}_w)$ is a rep of G

If $V \in \text{Irrep}(G)$:

$$\begin{aligned} & \text{multiplicity of } V \text{ in } L^2(\mathcal{O}_w) \\ &= \dim(H\text{-invariants} \subset V) \end{aligned}$$

e.g. $G = \text{SU}(2)$: $\text{Irrep}(\text{SU}(2)) = \{V_n \mid n = 0, 1, 2, \dots\}$

$$\mathcal{O}_w \cong S^2 \quad (w \neq 0), \quad L^2(\mathcal{O}_w) \cong \bigoplus_{n=0}^{\infty} V_{2n}.$$

Fuzzy sphere

Pick $V_n \in \text{Irrep}(su(2))$

$$V_n \cong \mathbb{C}^N, N=n+1$$

$$A = M_N(\mathbb{C})$$

Bimodule $\mathcal{H}_n = M_N(\mathbb{C})$ (or $M_N(\mathbb{C}) \otimes \mathbb{C}^2$)

$$\cong V_n \otimes V_n^* \quad \text{using adjoint action}$$

$$\cong V_0 \oplus V_2 \oplus V_4 \oplus \dots \oplus V_{2n}.$$

$$\mathcal{H}_n \subset \mathcal{H}_{n+1} \subset \dots \subset L^2(S^2)$$

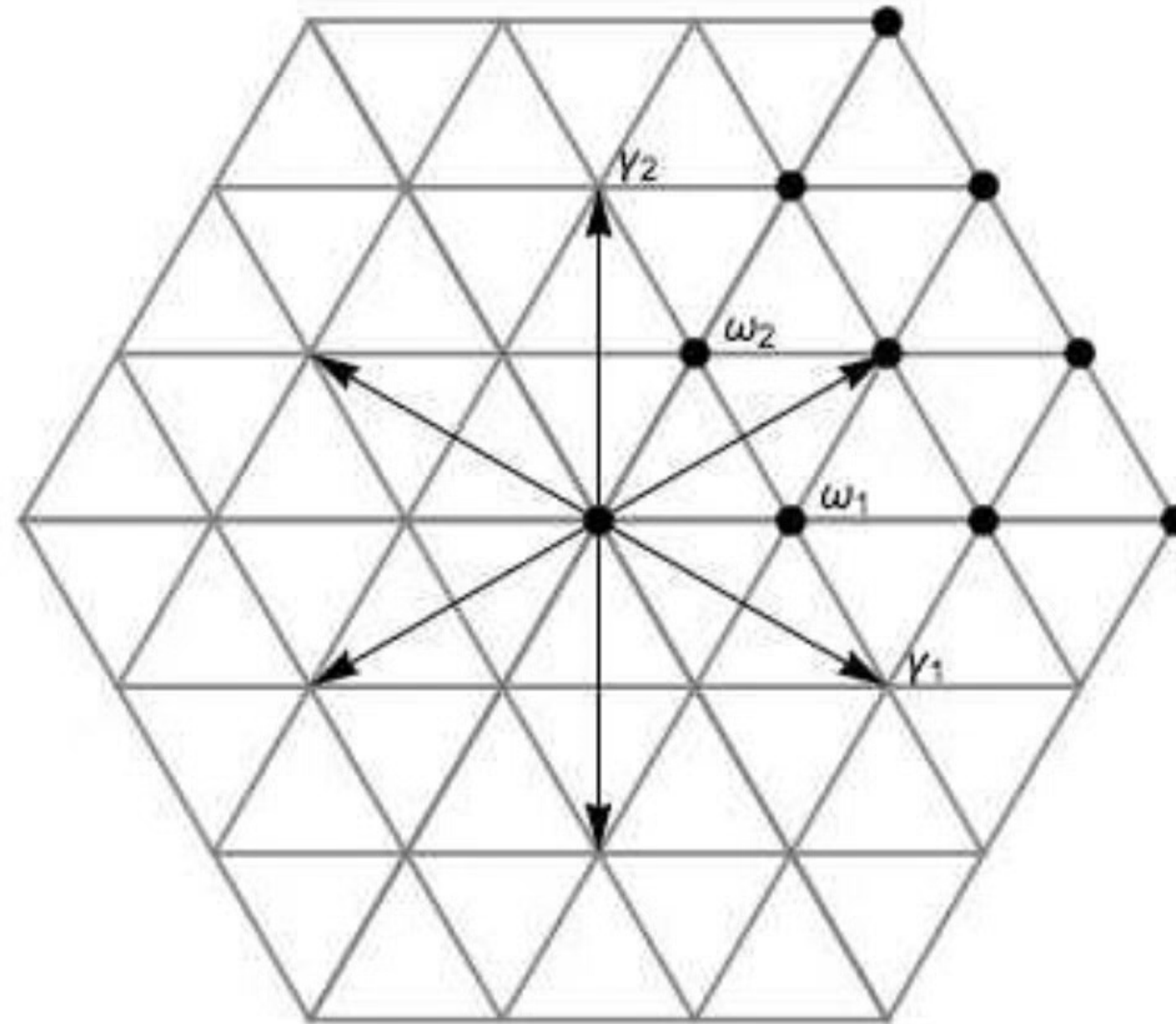
Fuzzy coadjoint orbit

$L^2(\mathcal{O}_w)$ is approximated by $V_w \otimes V_w^*$

In fact, $V_w \otimes V_w^* \subset V_{w'} \otimes V_{w'}^* \subset L^2(\mathcal{O}_{w'})$

if $w' - w \geq 0$

$SU(3)$ weight diagram



$$\omega_1 = (1, 0)$$

$$\omega_2 = (0, 1)$$

Algebraically integral elements (triangular lattice), dominant integral elements (black dots), and fundamental weights for $sl(3, \mathbb{C})$

Decomposition of tensor product of [5,3] and [3,5] in A2

Below you find the decomposition of the tensor product of the irreducible A2-modules with highest weights [5,3] and [3,5] into its irreducible factors; the result was computed by LiE using Klimyk's formula. Each term represents a different highest weight of an irreducible module occurring in the decomposition, prefixed by its multiplicity of occurrence. The tensor product has dimension 14400.

1X[8, 8] +1X[9, 6] +1X[6, 9] +1X[10, 4] +2X[7, 7] +1X[4,10] +
1X[11, 2] +2X[8, 5] +2X[5, 8] +1X[2,11] +2X[9, 3] +3X[6, 6] +
2X[3, 9] +1X[10, 1] +3X[7, 4] +3X[4, 7] +1X[1,10] +2X[8, 2] +
4X[5, 5] +2X[2, 8] +1X[9, 0] +3X[6, 3] +3X[3, 6] +1X[0, 9] +
2X[7, 1] +4X[4, 4] +2X[1, 7] +3X[5, 2] +3X[2, 5] +1X[6, 0] +
4X[3, 3] +1X[0, 6] +2X[4, 1] +2X[1, 4] +3X[2, 2] +1X[3, 0] +
1X[0, 3] +2X[1, 1] +1X[0, 0]

agree with $L^2(\mathcal{O}_W)$

Computation time 0.00 sec.

If you like, you may look at the [implementation](#) that was used (function *simp_tensor_irr* on page 2), which also involves routines for the [traversal of Weyl group orbits](#).

You may go back and try another example.

$$\mathcal{O}_W \cong \text{SU}(3) / \text{max. torus}$$

Euclidean group $E(2)$

Coadjoint orbits: $(p_1, p_2, \omega) \in L^*$

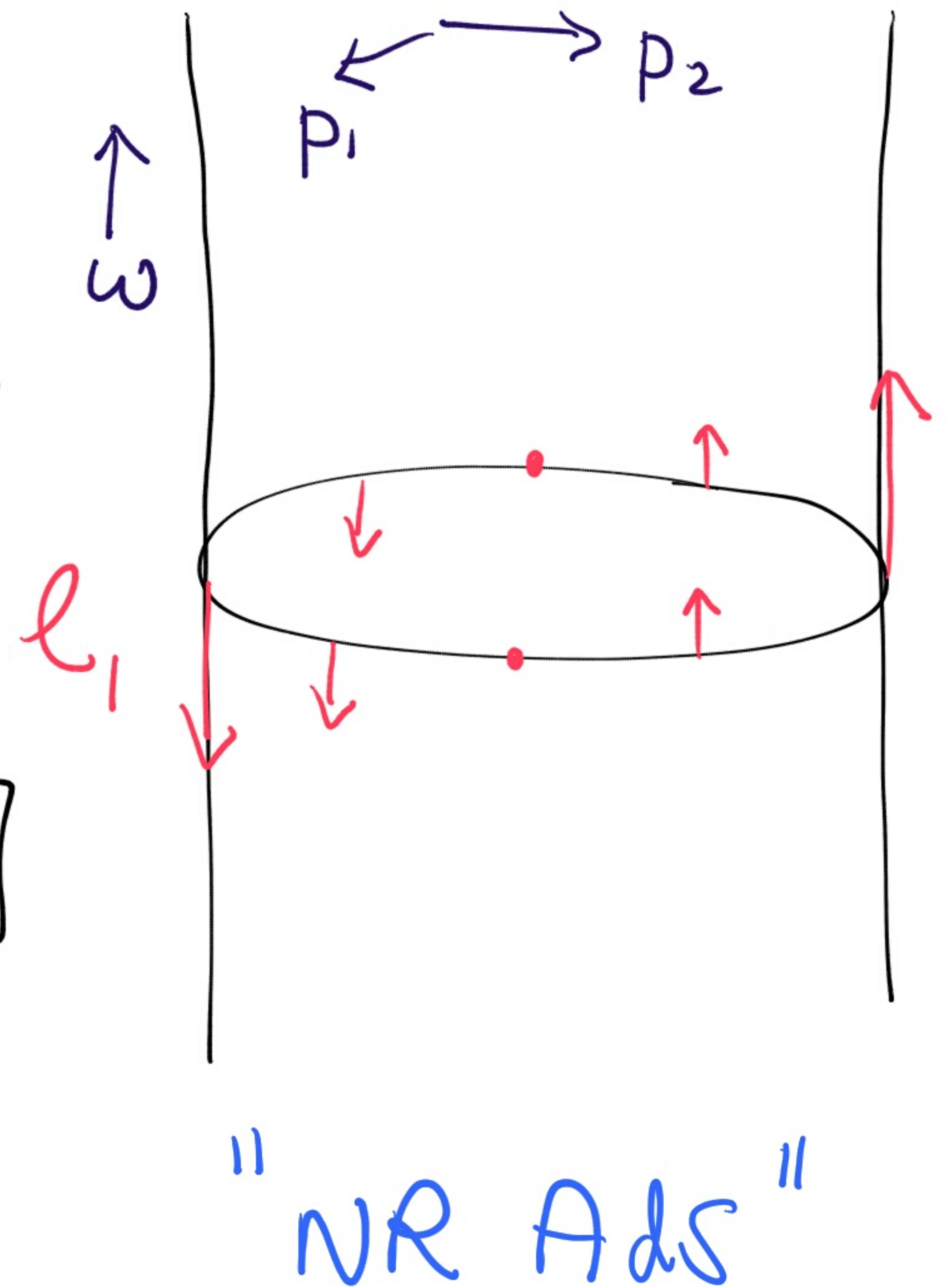
$$p_1^2 + p_2^2 = m^2 \neq 0$$

Laplacian $\Delta_{\mathcal{O}} = -\frac{\partial^2}{\partial \omega^2}$. Symplectic.

Fuzzy cylinder: $\mathcal{H}_m = V_m \otimes V_m^*$

$$\Delta_m \psi = -[l_1, [l_1, \psi]] - [l_2, [l_2, \psi]]$$

Spectrum of $\Delta_m = [0, 4m^2]$.



NC covering spaces

$$A \subset B \quad \text{* - algebras}$$

Define $G \subset A$ central unitaries (abelian group)

G acts on B : $b \mapsto gb g^{-1}$ deck transformations

If $\mathcal{H}$ is a bimodule for B , then

$$\mathcal{H} = \bigoplus_{\chi} \mathcal{H}_{\chi} \quad \chi: G \rightarrow U(1)$$

$\mathcal{H}_{\chi}$ are bimodules for A

Order Two coverings

M - manifold

s_1, s_2 spin structures

$$\leadsto s_1 - s_2 \in H_1(M, \mathbb{Z}_2)$$

$$\pi_1(M) \longrightarrow H_1(M; \mathbb{Z}_2)$$

$\mathbb{Z}_2$ -Hurewicz hom

determines covering $p: \hat{M} \rightarrow M$

with deck transformations $G = H_1(M; \mathbb{Z}_2)$

If $\hat{M}$ is compact, spin, $\mathcal{H} = L^2(\hat{M}, \text{spinors})$

then $\mathcal{H}_\chi \cong L^2(M, \text{spinors})$ with all spin structures

Fuzzy torus

$$B = M_{4N}(\mathbb{C}) = \langle C, S \rangle, \quad J = *$$

$$\text{with } CS = e^{2\pi i/4N} SC \quad \text{clock, shift}$$

$$C^{4N} = S^{4N} = 1$$

$$A = \langle U, V \rangle \text{ with } U = C^2, \quad V = S^2$$

$$\Rightarrow UV = e^{2\pi i/N} VU$$

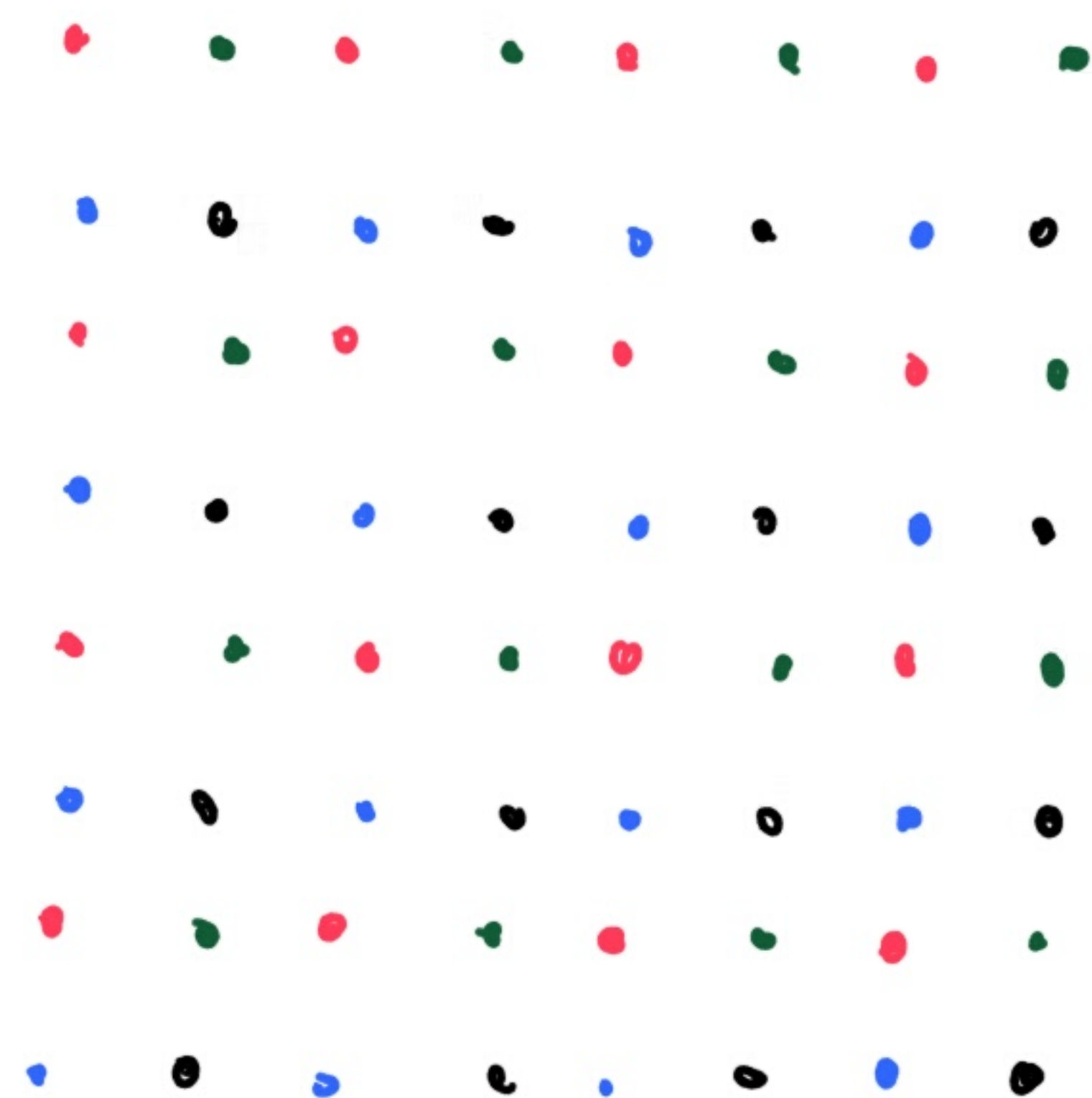
$$G = \langle U^N, V^N \rangle \cong \mathbb{Z}_2 \times \mathbb{Z}_2.$$

Fuzzy torus bimodules

$\mathcal{H} = M_{4N}(\mathbb{C}) \otimes \mathbb{C}^4$ has a Dirac operator
(JWB + James Gaunt)

$$\mathcal{H}_x \cong (M_N(\mathbb{C}) \oplus M_N(\mathbb{C}) \oplus M_N(\mathbb{C}) \oplus M_N(\mathbb{C})) \otimes \mathbb{C}^4$$

$N=4$
picture
of momentum
lattice



eigenvalues of D^2

$$[k]^2 + [l]^2$$

$$[k] = \frac{q^{k/2} - q^{-k/2}}{q^{1/2} - q^{-1/2}}$$